

Today, we will address the last (for now...) challenge associated with quasistatic robust simulation of non-linear materials \Rightarrow the positive-definiteness of the stiffness matrix K (or lack thereof)

Remember: We are using Newton's method to find the minimum-energy configuration for the function $E(\underline{x})$. We do so by iterating

$\underline{x}_{(0)} \leftarrow$ init. guess
for $k=0, 1, \dots$

$$\left[\begin{array}{l} \text{Solve } - \left. \frac{\partial f_{el}}{\partial \underline{x}} \right|_{\underline{x}_{(k)}} \cdot \delta \underline{x} = f_{el}(\underline{x}_{(k)}) \quad (1) \\ \text{Update } \underline{x}_{(k+1)} = \underline{x}_{(k)} + \delta \underline{x} \end{array} \right.$$

where eq (1) can also be seen as

$$K(\underline{x}_{(k)}) \delta \underline{x} = f_{el}(\underline{x}_{(k)}) \quad (2)$$

or even

$$\left. \frac{\partial^2 E}{\partial \underline{x}^2} \right|_{\underline{x}_{(k)}} \cdot \delta \underline{x} = - \left. \frac{\partial E}{\partial \underline{x}} \right|_{\underline{x}_{(k)}} \quad (3)$$

Our problem is that the matrix $K(\underline{x})$ is not guaranteed to be positive definite for any \underline{x}

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Quiz:

⇒ What would be an example of a configuration \underline{x} where $K(\underline{x})$ would be expected to be indefinite

⇒ Would you expect $K(\underline{x}_{(k)})$ to be positive-definite for a "late" iteration, where $\underline{x}_{(k)}$ is almost converged and why? (or why not?)

The general philosophy behind our remedy:

⇒ Replace the stiffness matrix $K(\underline{x}_{(k)})$ with a modified version $\hat{K}(\underline{x}_{(k)})$ which is designed to be positive-definite (CG can be used!)

⇒ The method becomes modified Newton, iterating

$$\hat{K}(\underline{x}_{(k)}) \delta \underline{x} = \underline{f}_{el}(\underline{x}_{(k)})$$

But why is this a reasonable remedy?

⇒ It depends on the nature of the modification (\hat{K}) but we can say the following:

→ For the specific method we will see, practice shows that modified Newton is very robust, especially when combined with inversion-robust fixes (truncating singular values to small positive).

→ Once the modified iteration is close enough to convergence (enough for $K(x)$ to be s.p.d) we can show that

$$\rho \left(I - [\hat{K}(x_{(k)})]^{-1} K(x_{(k)}) \right) < 1$$

↑ spectral radius

then the modified Newton process is guaranteed to converge to the actual minimum.

(i.e. if we are almost converged, a small modification does not compromise the ability to converge all the way)

(In-class discussion: 2D nonzero-rest length spring

$$E(\vec{x}) = \frac{k}{2} (\|\vec{x} - \vec{x}_0\| - l_0)^2$$

Outline of the modification

⇒ Split the stiffness matrix to the contribution of individual elements

$$K(\underline{x}) = \sum_{e: \text{element}} K^e(\underline{x})$$

Each K^e corresponds to the contribution of just the e -th element to the energy, i.e.:

$$\text{if } E(\underline{x}) = \sum_e E^e(\underline{x})$$

$$\text{then } K^e := - \frac{\partial E^e(\underline{x})}{\partial \underline{x}}$$

if (i, j, k, l) are the vertex indices, then only the blocks of K^e corresponding to row/column blocks $i/j/k/l$ will be non-zero

$$K_e = \begin{bmatrix} \circ & \circ & \circ & \circ & \circ & \circ & \circ & \circ \\ \circ & X & & X & X & X & & \\ \circ & & \circ & & & & & \\ \circ & X & & X & X & X & & \\ \circ & & & X & X & X & & \\ \circ & X & & X & X & X & & \\ \circ & X & & X & X & X & & \\ \circ & X & & X & X & X & & \\ \uparrow & \uparrow & \uparrow & \uparrow & & & & \\ i & j & k & l & & & & \end{bmatrix}$$

Each "X" is a 3x3 matrix, so, effectively K_e is a 12x12 matrix padded by zeroes to become full-size

⇒ The prescription for the remedy, then, becomes:

→ "Project" each elemental stiffness matrix to its positive definite part

$$K^e = Q \Lambda Q^T$$

↑ only 12 of the λ_i 's can be nonzero

(in fact, just 9 of them maximum ... why?)

If we define $(\hat{\Lambda})_{ii} = \begin{cases} (\Lambda)_{ii} & , \text{ if } \Lambda_{ii} \geq 0 \\ 0 & , \text{ if } \Lambda_{ii} < 0 \end{cases}$

and $\hat{K}^e := Q \hat{\Lambda} Q^T$

→ Construct the "global" modified stiffness by assembly:

$$\hat{K} := \sum_e \hat{K}^e$$

(some textbooks use an "assembly" operator A to write $\hat{K} := A K^e$)

where K^e are 12×12 matrices)

If we were prepared to pay the cost of a 12×12 eigenanalysis (9×9 , with some effort...) this concludes our fix! (in some cases, we would have no other option).

There is, however, a shortcut to computing exactly the same result, with much reduced cost in the case of isotropic materials (those where $\psi = \psi(\Sigma)$)

Implementation

Remember the matrix-free computation of the differential $\delta f[x; \delta x]$

add Force Differential

foreach element $e = (i, j, k, l)$

$D_s^e, F^e \leftarrow$ as before

$\delta D_s^e, \delta F^e \leftarrow$ as before

$$\delta P^e = \delta P[F; \delta F]$$

$$\delta H^e \leftarrow -Vole \cdot \delta P^e (D_m^e)^{-1}$$

$$\delta H^e \rightarrow \delta f_{ijkl} \text{ as before}$$

Definiteness fix

Replace $\delta P[F; \delta F]$ with modified

$$\hat{\delta P}[F; \delta F]$$

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And, here is the prescription for the modified "stress differential"

(Proofs: Teran et al '05 "Robust Quasistatics and Flesh Simulation", SCA '05
Stomakhin et al '12 "Energetically Consistent Invertible Elasticity", SCA '12)

Definitions

$$\text{vec} \begin{bmatrix} \begin{pmatrix} a_{11} & a_{12} & a_{13} \\ a_{21} & a_{22} & a_{23} \\ a_{31} & a_{32} & a_{33} \end{pmatrix} \end{bmatrix} = \begin{pmatrix} a_{11} \\ a_{22} \\ a_{33} \\ a_{12} \\ a_{21} \\ a_{13} \\ a_{31} \\ a_{23} \\ a_{32} \end{pmatrix}$$

$$\text{mat} \begin{bmatrix} \begin{pmatrix} a_1 \\ a_2 \\ a_3 \\ a_4 \\ a_5 \\ a_6 \\ a_7 \\ a_8 \\ a_9 \end{pmatrix} \end{bmatrix} = \begin{pmatrix} a_1 & a_4 & a_6 \\ a_5 & a_2 & a_8 \\ a_7 & a_9 & a_3 \end{pmatrix}$$

We can show that $\delta I = U \text{mat} [T \cdot \text{vec} (U^T \delta F V)] V^T$
($F = U \Sigma V^T$)

Where $T = \begin{bmatrix} A_{(3 \times 3)} & & & \\ & B_{12}_{(2 \times 2)} & & \\ & & B_{13}_{(2 \times 2)} & \\ & & & B_{23}_{(2 \times 2)} \end{bmatrix}$

where $A_{ij} = \frac{\partial^2 \psi}{\partial \sigma_i \partial \sigma_j}$

$$B_{ij} = \frac{1}{2} \frac{\frac{\partial \psi}{\partial \sigma_i} - \frac{\partial \psi}{\partial \sigma_j}}{\sigma_i - \sigma_j} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} + \frac{1}{2} \frac{\frac{\partial \psi}{\partial \sigma_i} + \frac{\partial \psi}{\partial \sigma_j}}{\sigma_i + \sigma_j} \begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix}$$